

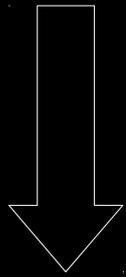
# New approach to the local strong parity violation in the quark-gluon plasma



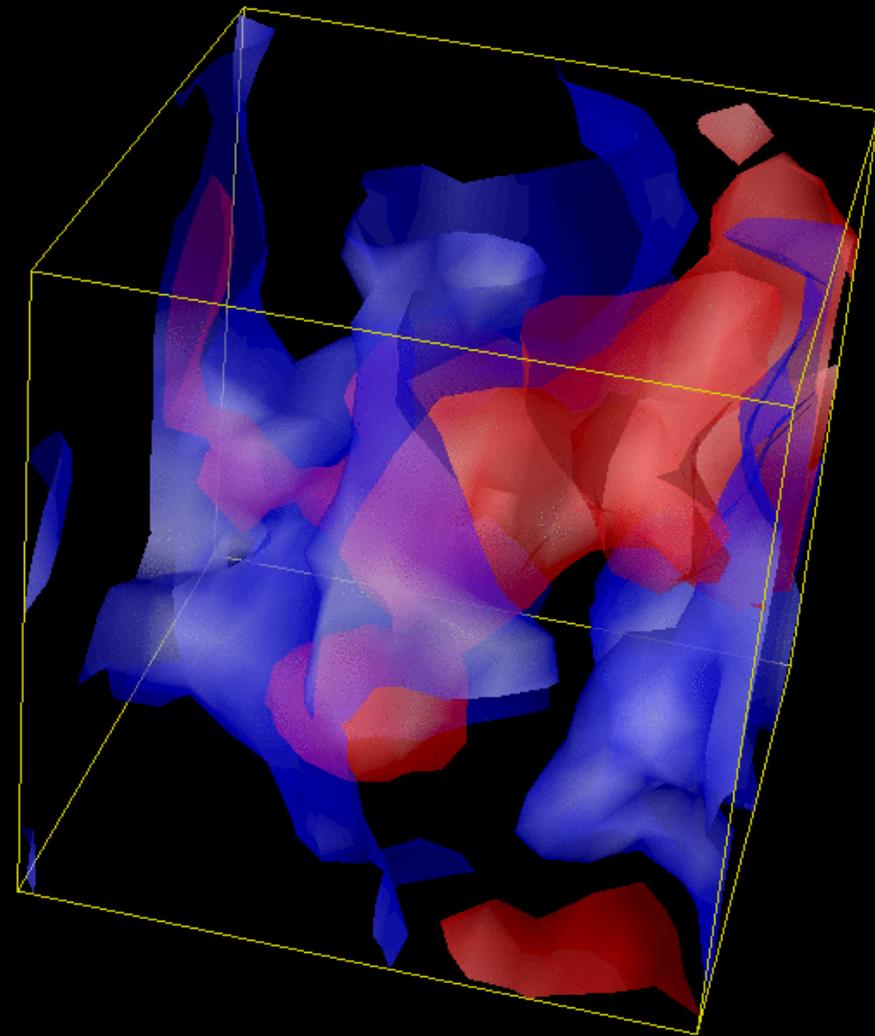
Tigran Kalaydzhyan

# QCD vacuum

$$G^{a\mu\nu} \tilde{G}_{\mu\nu}^a$$



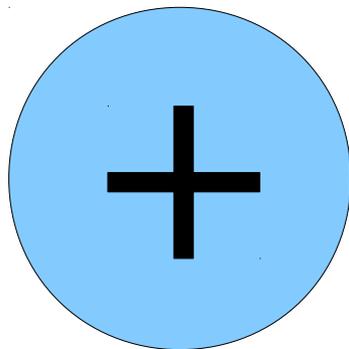
$$\rho_R \neq \rho_L$$



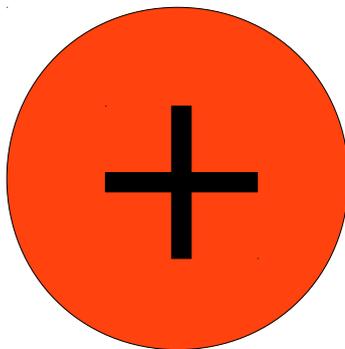
Positive topological  
charge density

Negative topological  
charge density

# (Naive) visible effects

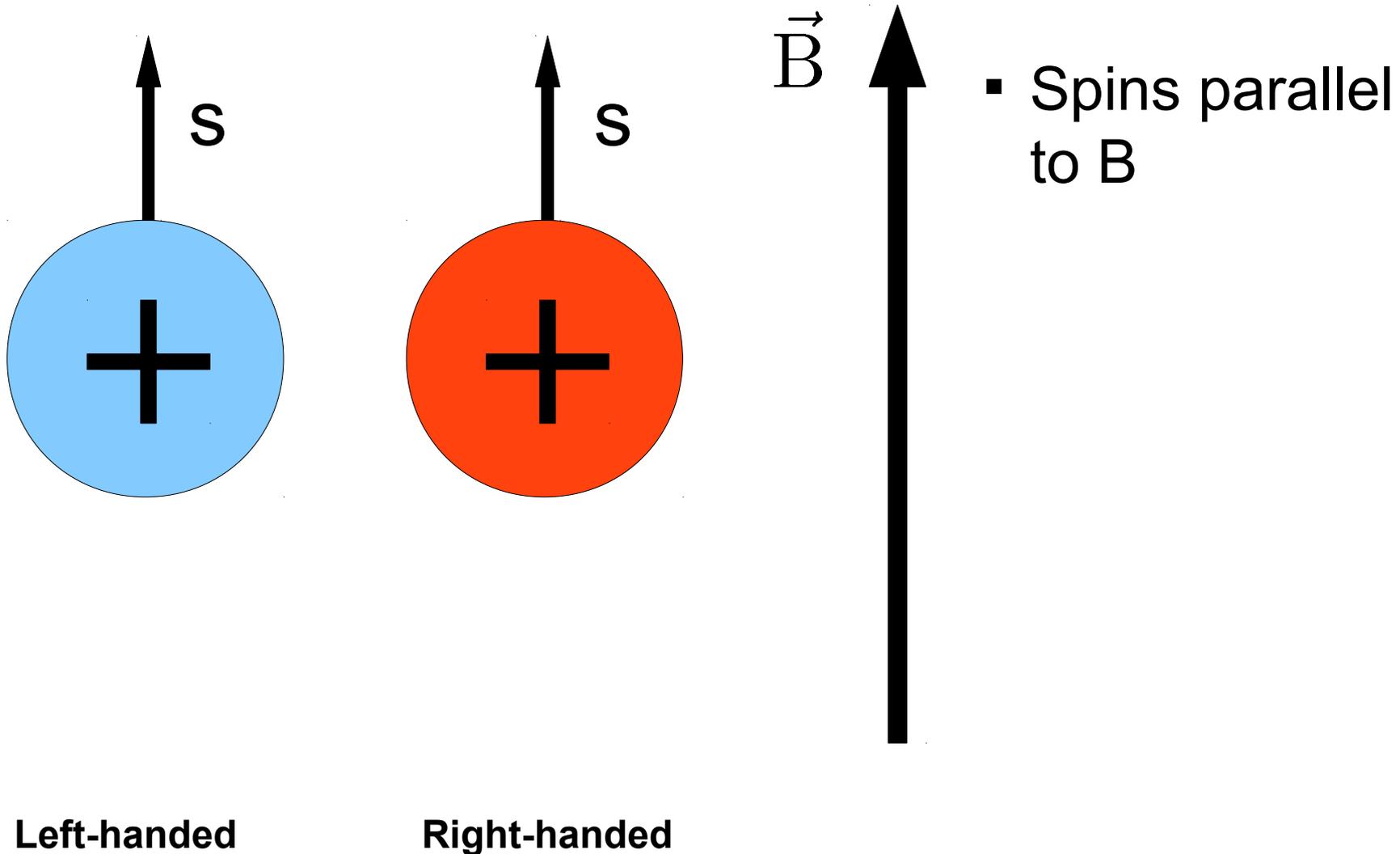


Left-handed

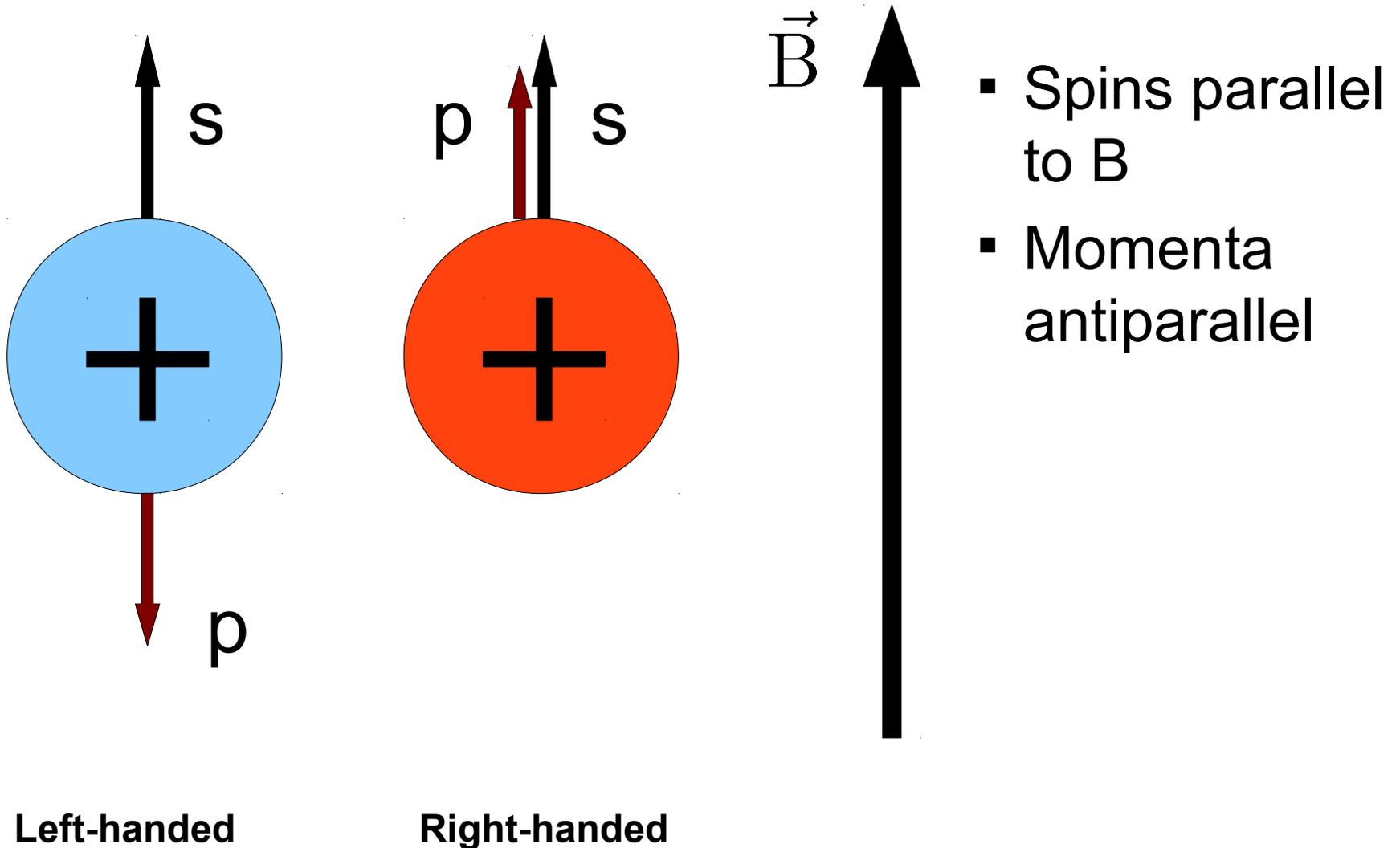


Right-handed

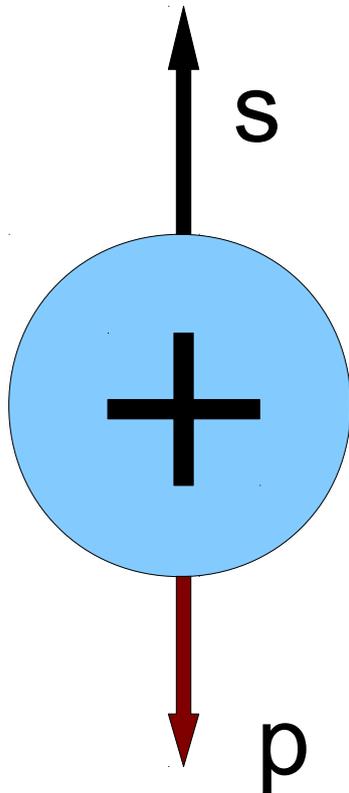
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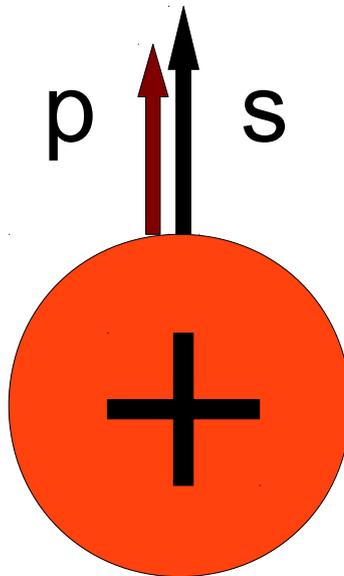
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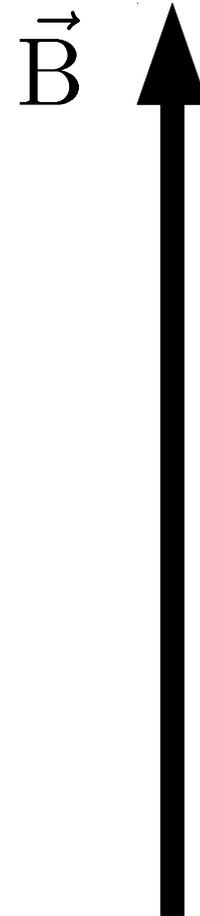
# (Naive) visible effects



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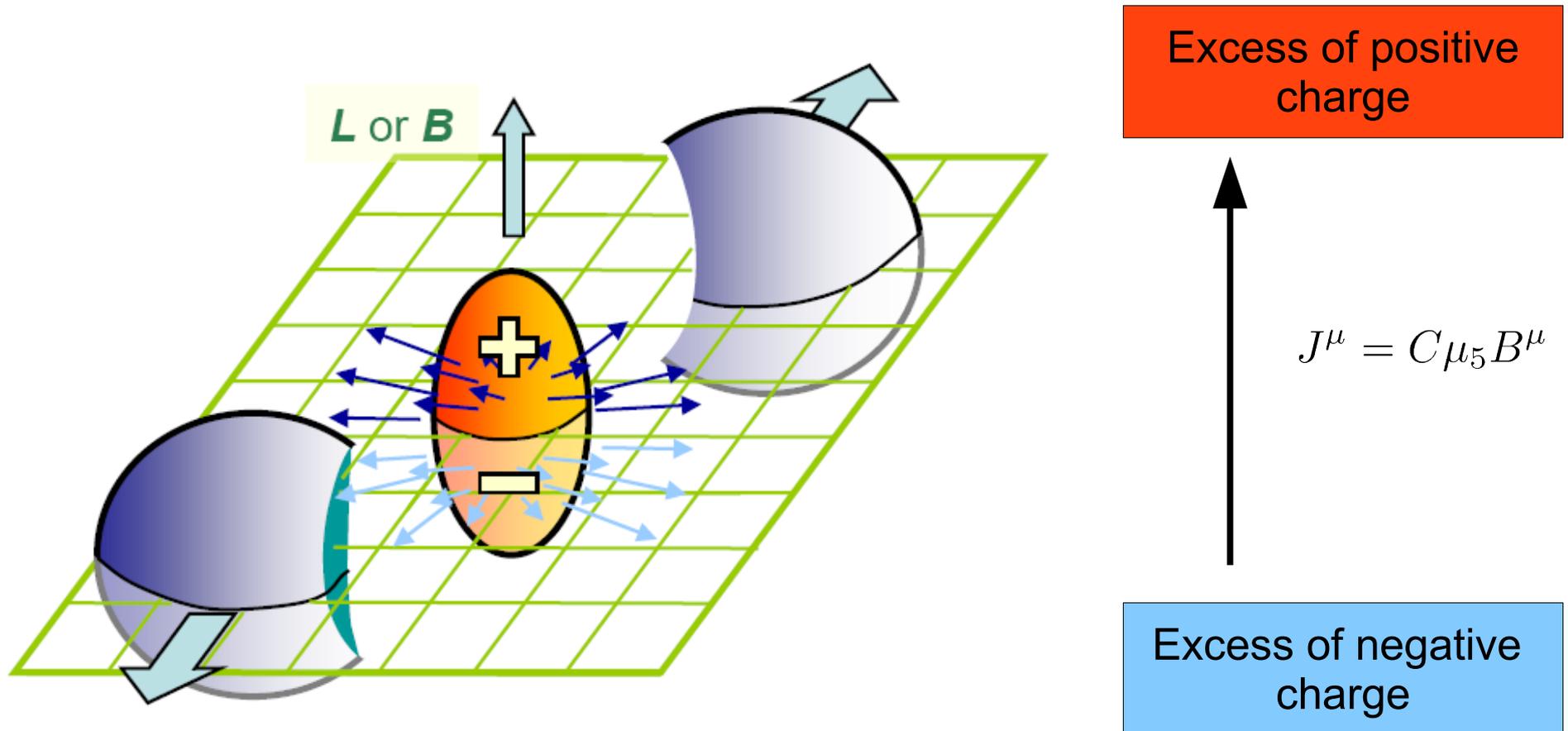


Right-handed



- Spins parallel to  $B$
- Momenta antiparallel
- If  $\rho_5 \equiv \rho_L - \rho_R \neq 0$  then we have a net electric current parallel to  $B$

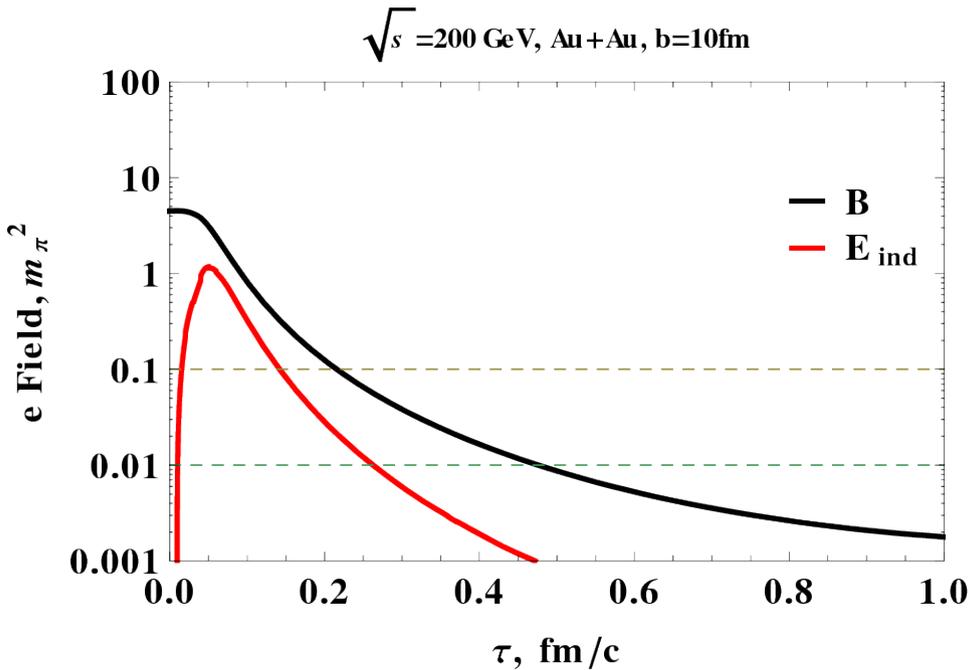
# Chiral Magnetic Effect



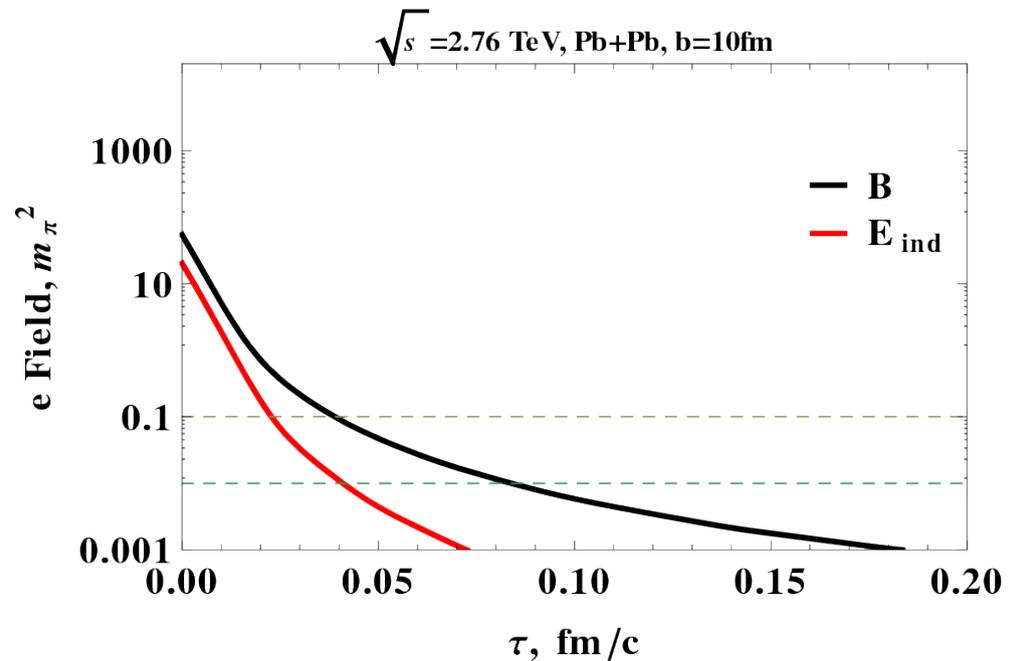
Fukushima, Kharzeev, McLerran, Warringa (2007)

For a local strong parity violation see e.g. 0909.1717 (STAR) and 1207.0900 (ALICE)

# Electromagnetic fields



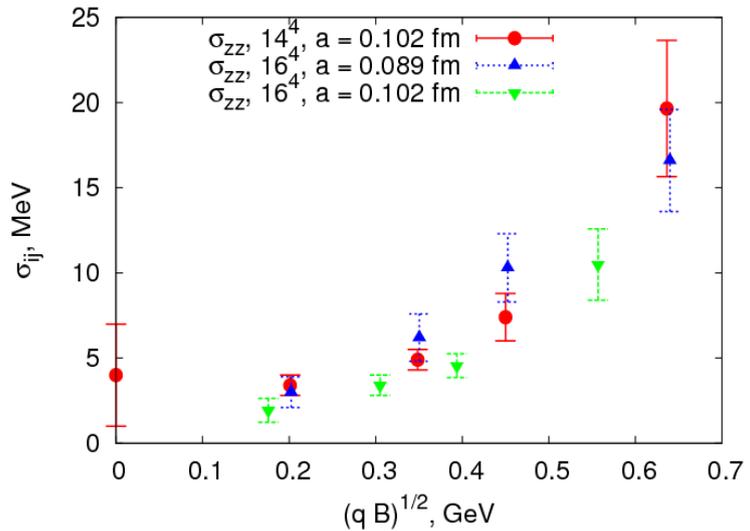
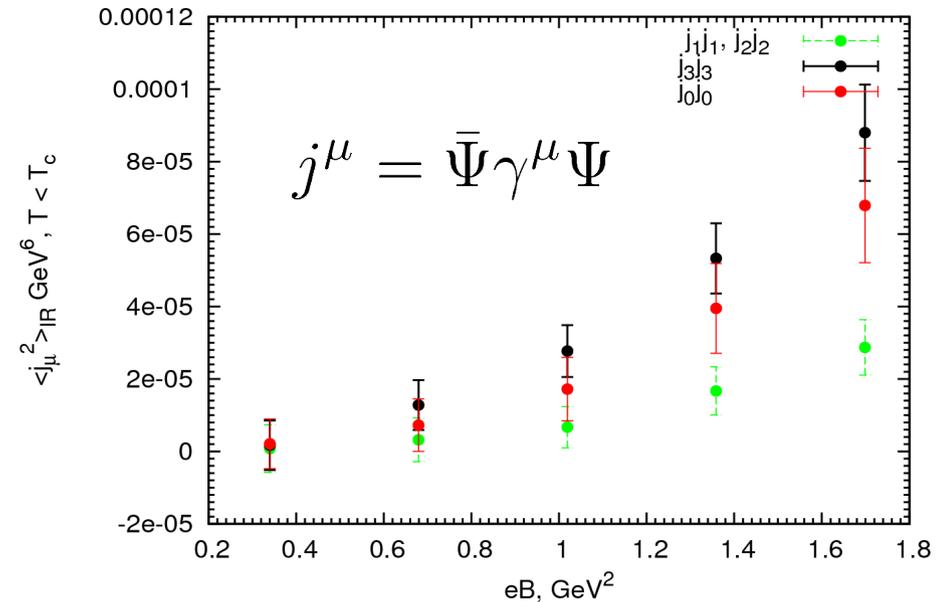
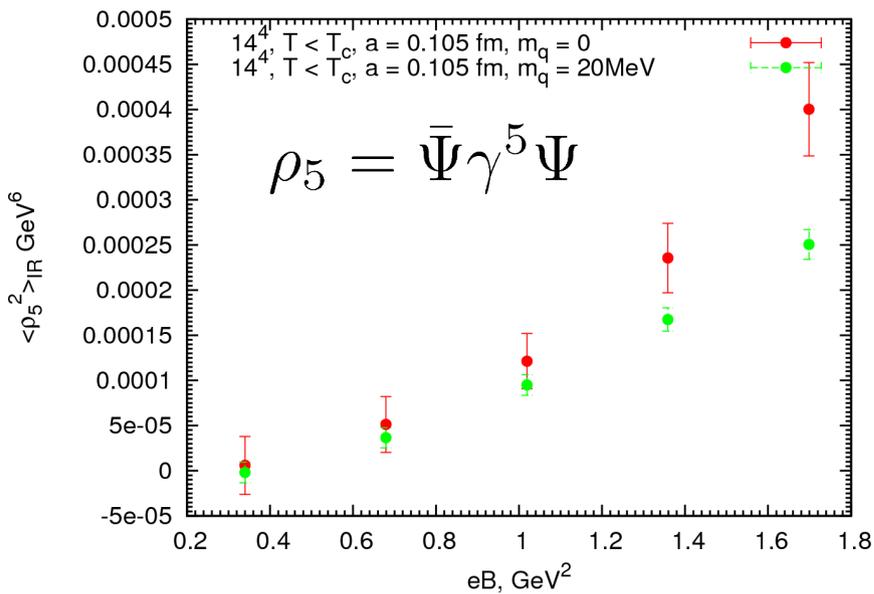
**RHIC**



**LHC**

**Huge electromagnetic fields, never observed before!**

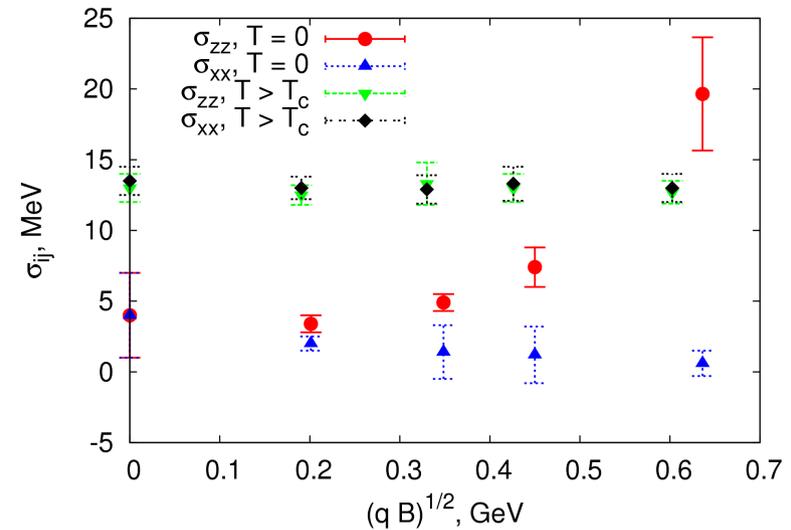
# Some numbers (lattice)



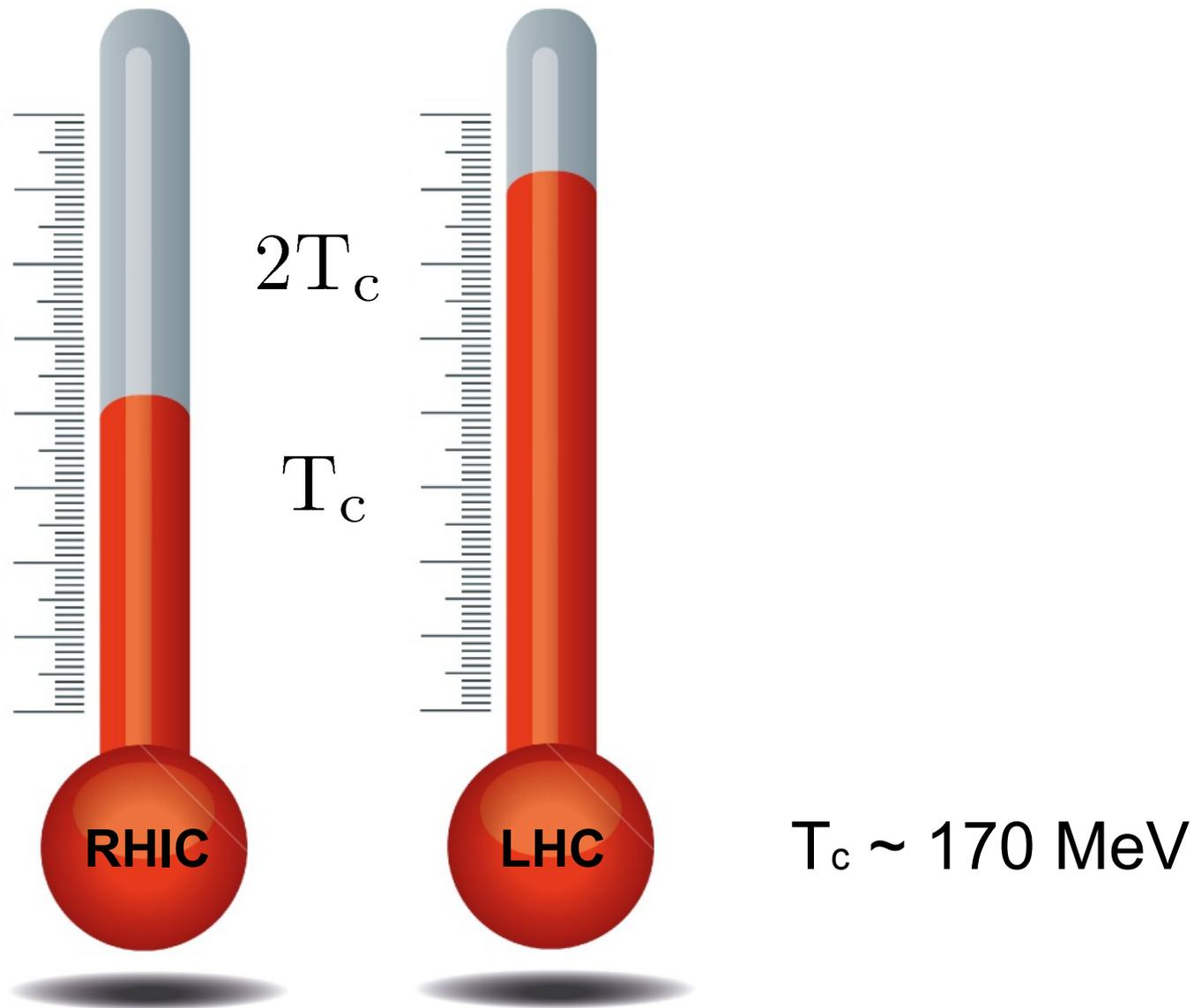
T.K., D. Kharzeev and



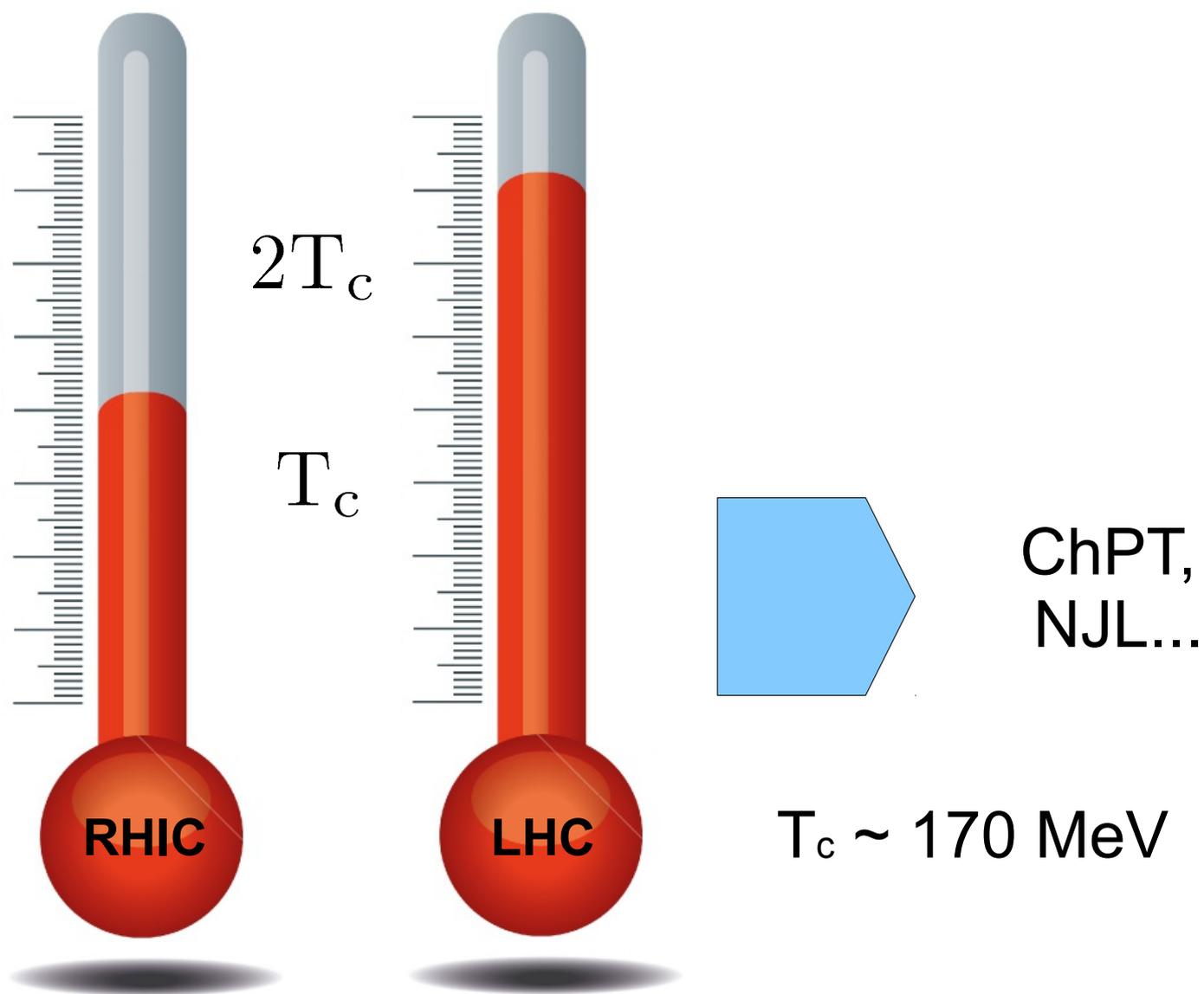
PRL 105 (2010) 132001  
 Phys.Atom.Nucl. 75, 488



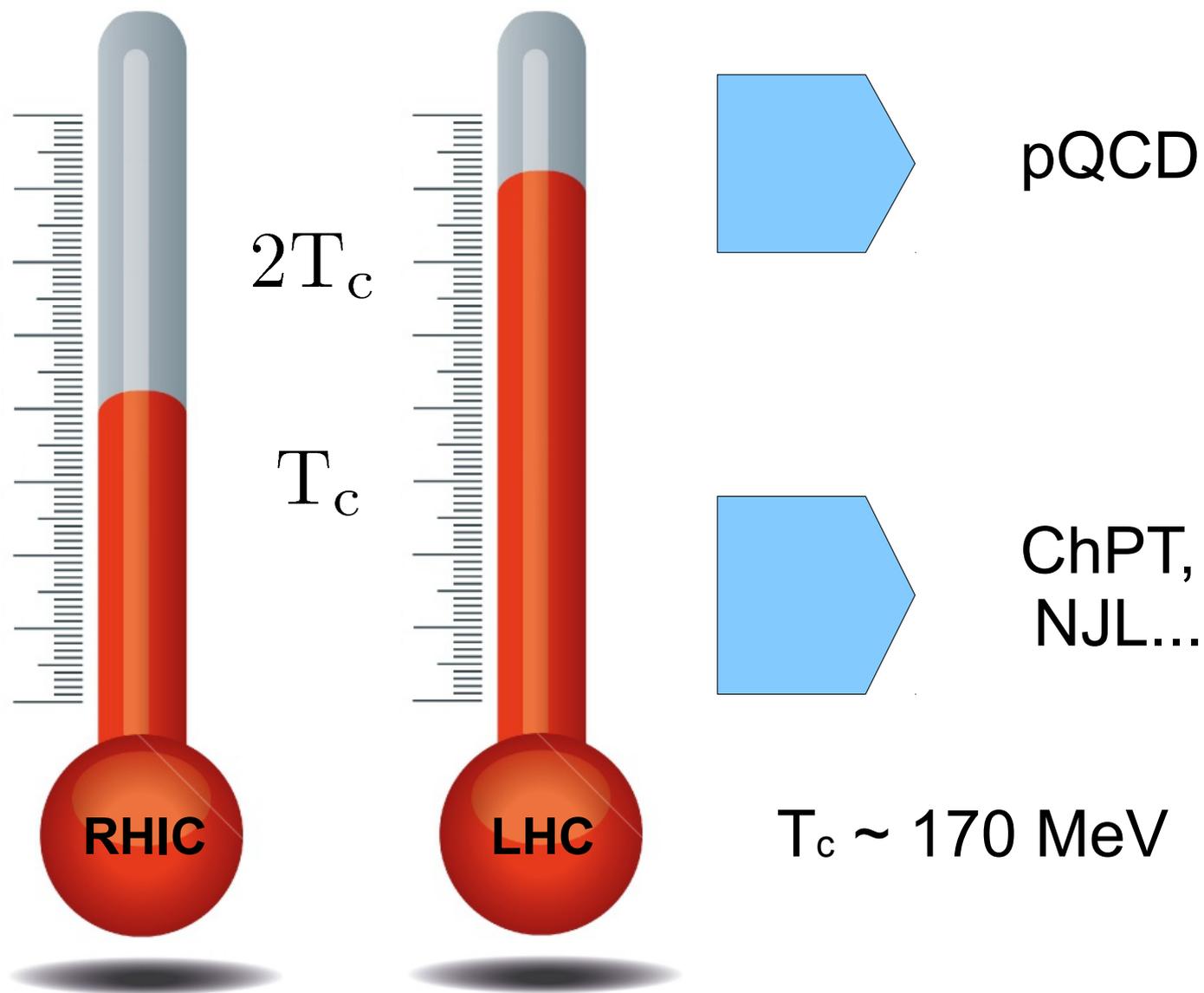
# Temperatures



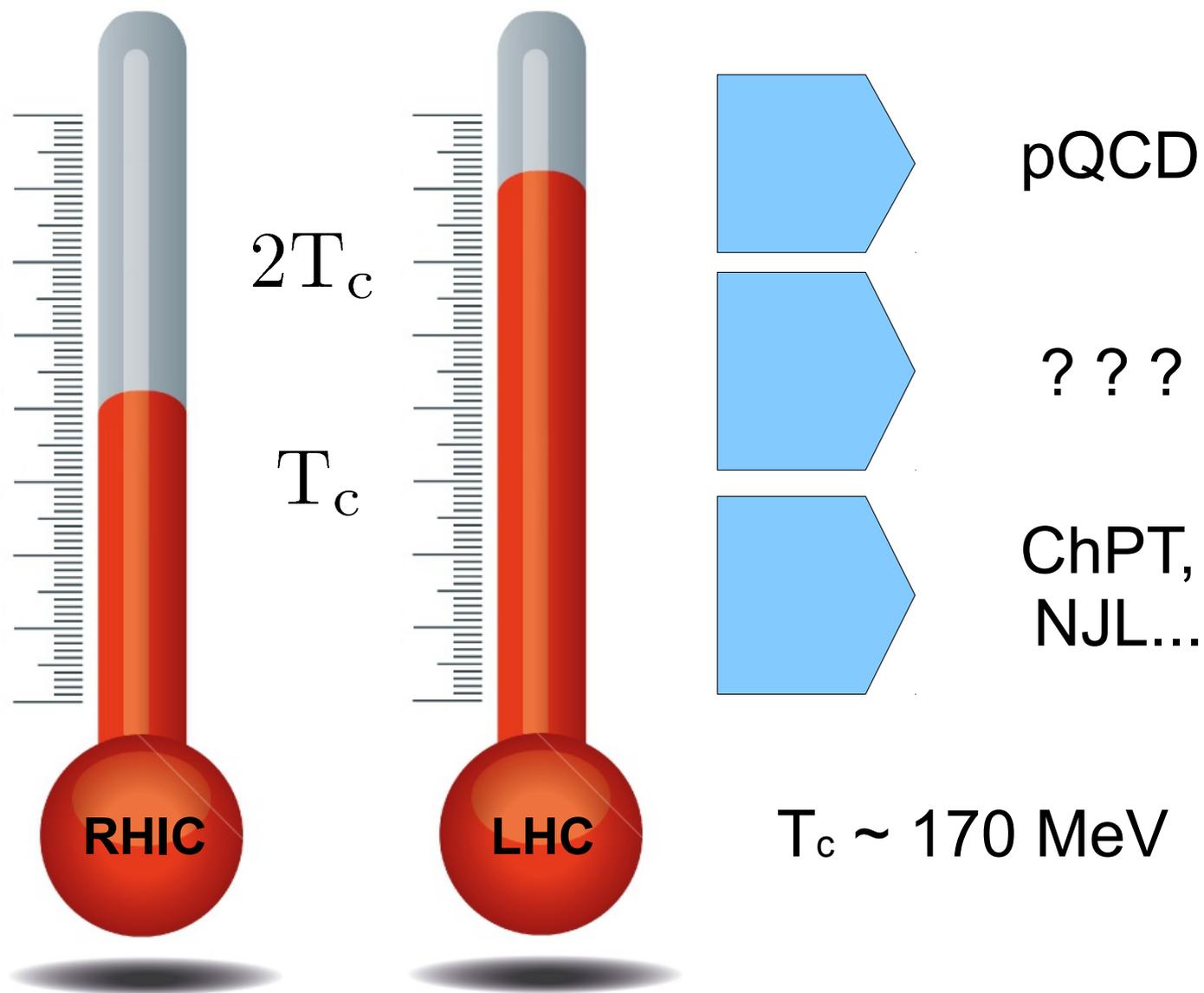
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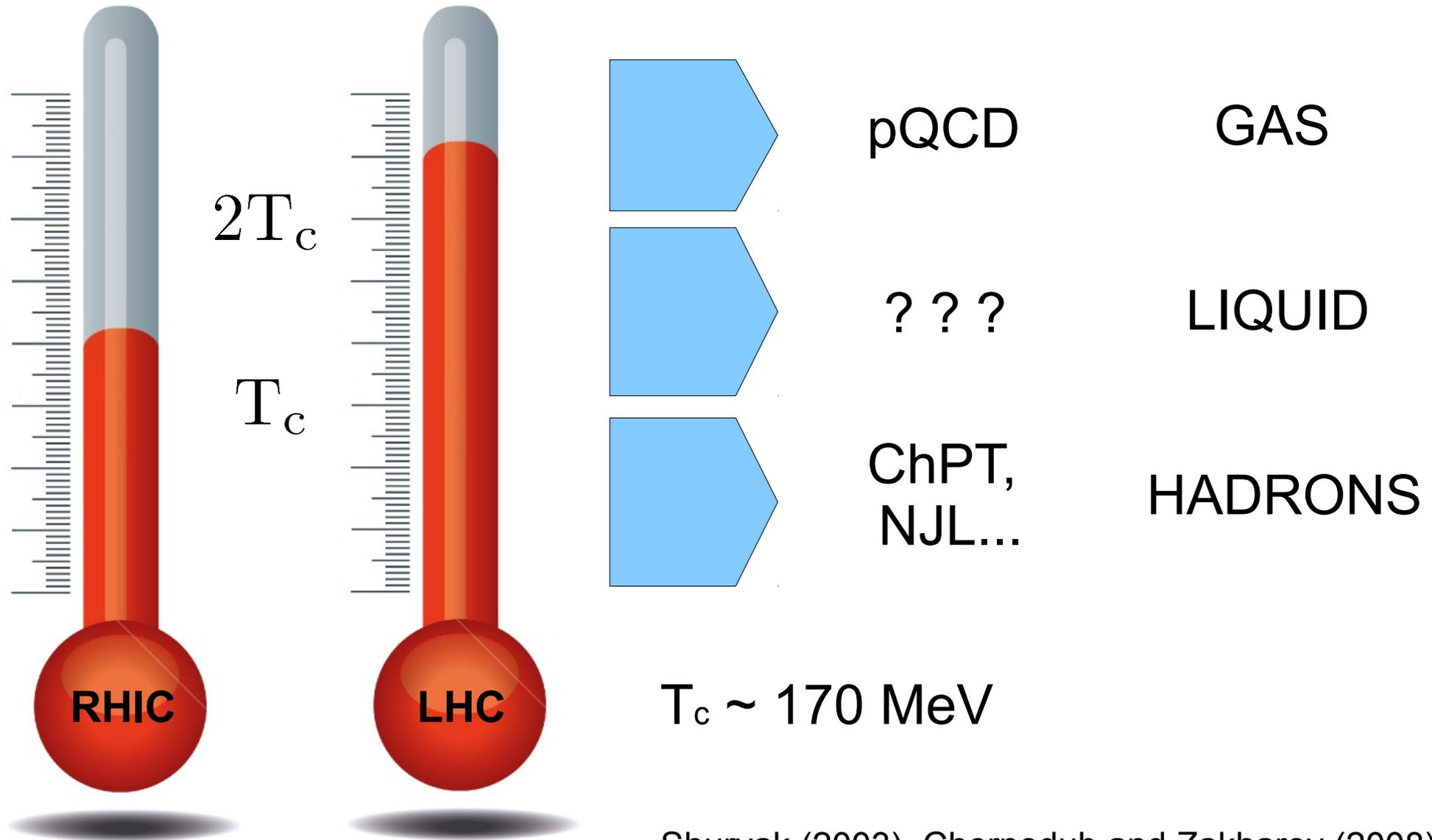
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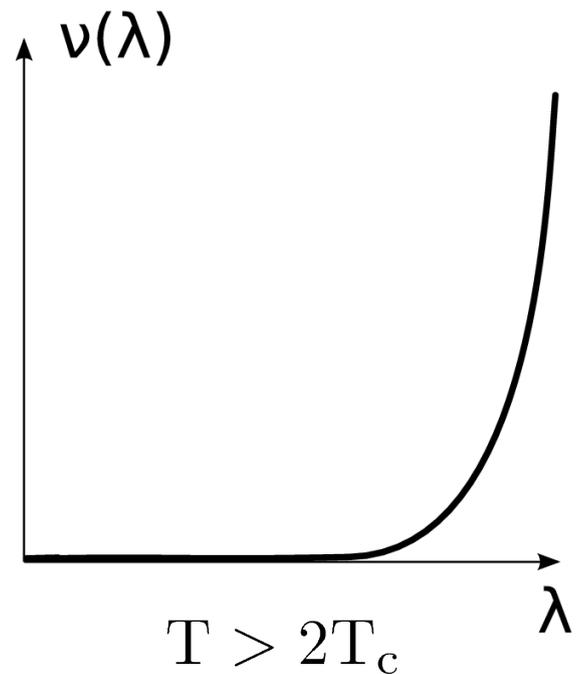
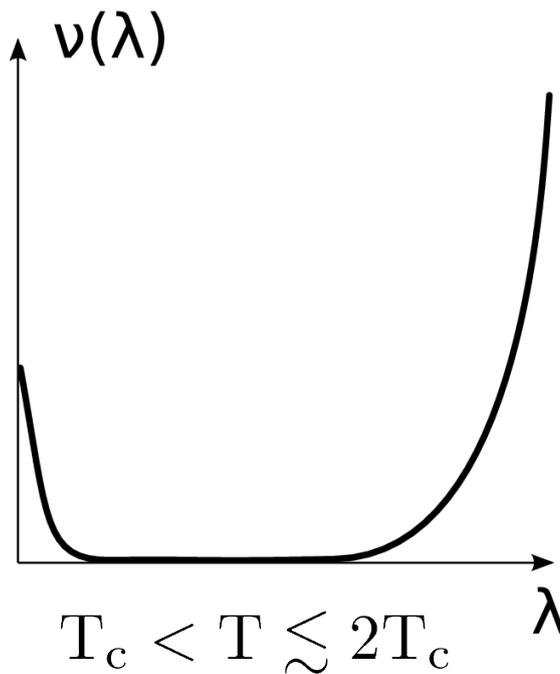
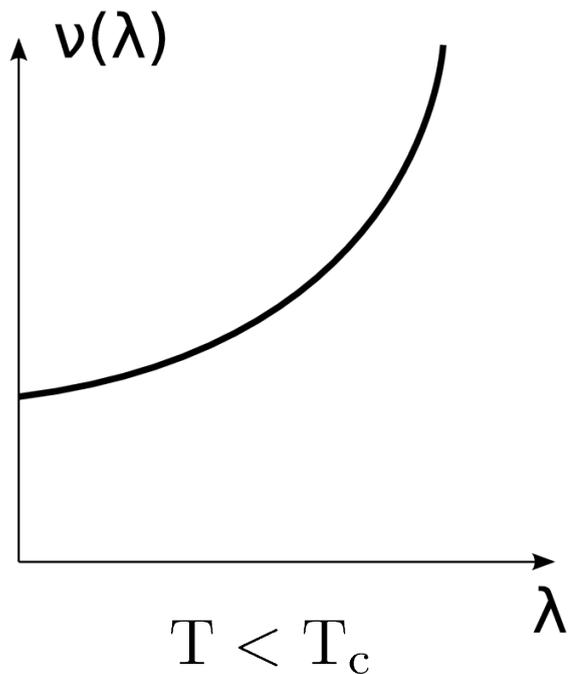
# Our task

- Build an effective model for QCD at  $T_c < T < 2 T_c$  (did we lose anything?)
- Find hydrodynamic equations corresponding to the effective Lagrangian
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- Extract phenomenological output for the heavy-ion collisions

# Insight from the lattice

- Spectrum of the Dirac operator

$$\hat{D}\psi_\lambda = \lambda\psi_\lambda$$

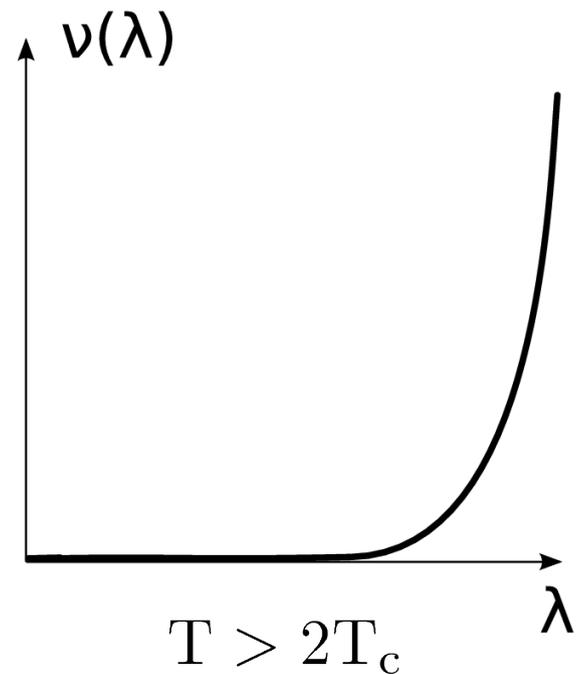
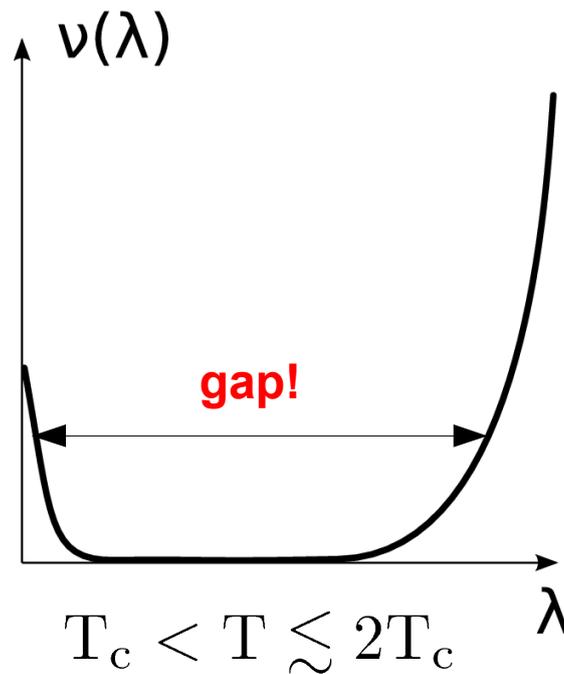
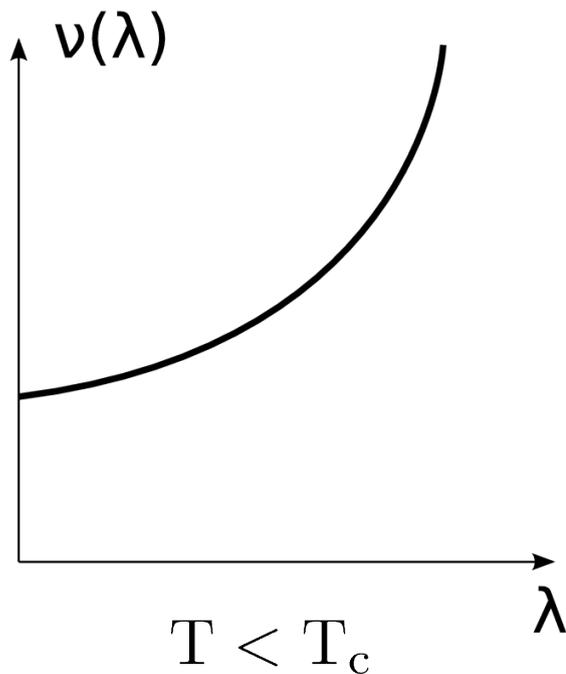


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# Insight from the lattice

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- Chiral properties are described by near-zero modes
- There are two separated parts of the spectrum at intermediate temperatures!

# Why "superfluidity" ?

AUGUST 15, 1941

PHYSICAL REVIEW

## Theory of the Superfluidity of Helium II

L. LANDAU

*Institute of Physical Problems, Academy of Sciences USSR, Moscow, USSR*

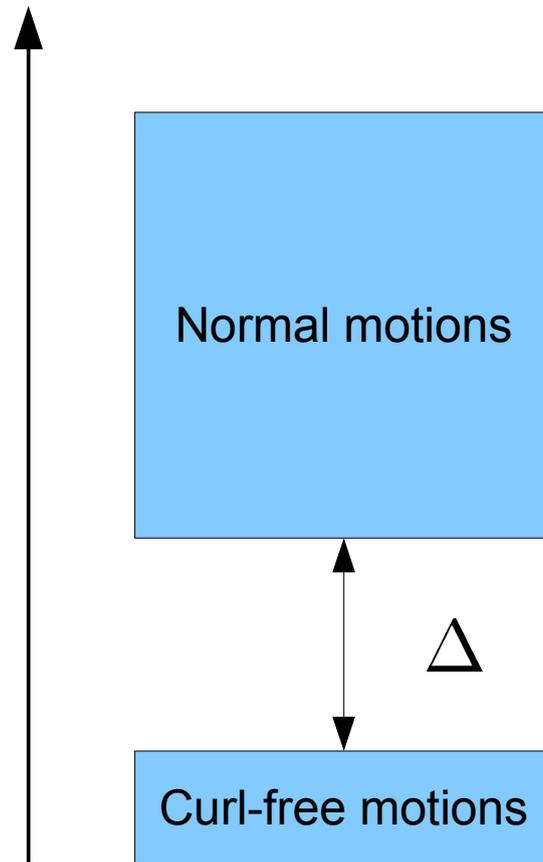
Therefore, between the lowest energy levels of vortex and potential motion there must be a certain energy interval  $\Delta$ .

The supposition that the normal level of potential motions lies lower than the beginning of the spectrum of vortex motions leads to the phenomenon of superfluidity.

One of these motions is "normal" and the other is "superfluid."

**We will not consider any spontaneously broken symmetry!**

Energy



# Bosonization

- Euclidean functional integral for  $SU(N_c) \times U_{em}(1)$  is given by

$$\int D\bar{\psi} D\psi \exp \left\{ - \int_V d^4x \bar{\psi} (\not{D} - im) \psi + \frac{1}{4} G^{a\mu\nu} G_{\mu\nu}^a + \frac{1}{4} F^{\mu\nu} F_{\mu\nu} \right\},$$

where we define the Dirac operator as

$$\not{D} = -i(\not{\partial} + \not{A} + g\not{G} + \gamma_5 \not{A}_5).$$

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- and the chiral limit  $m \rightarrow 0$

# Bosonization

- The **total effective Euclidean Lagrangian** reads as

$$\begin{aligned}\mathcal{L}_E^{(4)} &= \frac{1}{4} G^{a\mu\nu} G_{\mu\nu}^a + \frac{1}{4} F^{\mu\nu} F_{\mu\nu} - j^\mu A_\mu \\ &+ \frac{\Lambda^2 N_c}{4\pi^2} \partial^\mu \theta \partial_\mu \theta + \frac{g^2}{16\pi^2} \theta G^{a\mu\nu} \tilde{G}_{\mu\nu}^a + \frac{N_c}{8\pi^2} \theta F^{\mu\nu} \tilde{F}_{\mu\nu} \\ &+ \frac{N_c}{24\pi^2} \theta \square^2 \theta - \frac{N_c}{12\pi^2} (\partial^\mu \theta \partial_\mu \theta)^2\end{aligned}$$

So we get an axion-like field with decay constant  $f = \frac{2\Lambda}{\pi} \sqrt{N_c}$

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**Dynamical axion-like internal degree of freedom in QCD!**

# Interpretation of the scale $\Lambda$

- From the quartic Lagrangian at  $N_c = N_f = 1$  we get

$$\rho_5 = - \lim_{t_E \rightarrow 0} \frac{\delta \mathcal{L}_E^{(4)}}{\delta \mu_5} = \frac{1}{2} \left( \frac{\Lambda}{\pi} \right)^2 \mu_5 + \frac{1}{3\pi^2} \mu_5^3$$

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- Free quarks and a strong B-field:  $\Lambda = 2\sqrt{|eB|}$
- Dynamical fermions (1105.0385):  $\Lambda \simeq 3 \text{ GeV} \gg \Lambda_{QCD}$

**A „hidden“ scale!**

# One more remark

„Axionic“ part of the Lagrangian

$$\mathcal{L}_\theta = \frac{\Lambda^2 N_c}{4\pi^2} \partial^\mu \theta \partial_\mu \theta + \frac{N_c}{24\pi^2} \theta \square^2 \theta - \frac{N_c}{12\pi^2} (\partial^\mu \theta \partial_\mu \theta)^2 + \dots$$

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If magnetic field dominates over other scales, then we can make

the following redefinition:  $\theta \rightarrow \frac{\pi}{\sqrt{2N_c eB}} \theta$

$$\mathcal{L}_\theta \rightarrow \frac{1}{2} \partial^\mu \theta \partial_\mu \theta + \frac{1}{48 eB} \theta \square^2 \theta - \frac{\pi^2}{48 N_c (eB)^2} (\partial^\mu \theta \partial_\mu \theta)^2 + \dots$$

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In the limit  $B \rightarrow \infty$  bosonization becomes exact, which is an evidence of the  $(3+1) \rightarrow (1+1)$  reduction!

# Hydrodynamic equations

Considering EOM for the Minkowski effective Lagrangian and only the color-singlet states, we obtain:

$$\partial_\mu T^{\mu\nu} = F^{\nu\lambda} J_\lambda ,$$

$$\partial_\mu J^\mu = 0 ,$$

$$\partial_\mu J_5^\mu = C E^\mu B_\mu ,$$

$$u^\mu \partial_\mu \theta + \mu_5 = 0 ,$$

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Energy-momentum tensor  $\rightarrow \partial_\mu T^{\mu\nu} = F^{\nu\lambda} J_\lambda,$

$\partial_\mu J^\mu = 0,$

Axial current  $\rightarrow \partial_\mu J_5^\mu = C E^\mu B_\mu,$

4-velocity of the normal component  $\rightarrow u^\mu \partial_\mu \theta + \mu_5 = 0,$

Total electric current

Electromagnetic fields

Chiral anomaly coefficient

Josephson equation, defining The axial chemical potential through the bosonized low-lying modes.

**Similar to the superfluid dynamics!**

# Hydrodynamic equations

- Solving the hydrodynamic equations in the gradient expansion, we obtain the constitutive relations:

$$T^{\mu\nu} = (\epsilon + P) u^\mu u^\nu + P g^{\mu\nu} + f^2 \partial^\mu \theta \partial^\nu \theta + \tau^{\mu\nu},$$

$$J^\mu = \rho u^\mu + C \tilde{F}^{\mu\kappa} \partial_\kappa \theta + \nu^\mu,$$

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$$T^{\mu\nu} = (\epsilon + P) u^\mu u^\nu + P g^{\mu\nu} + f^2 \partial^\mu \theta \partial^\nu \theta + \tau^{\mu\nu},$$

Energy density

Pressure

$$J^\mu = \rho u^\mu + C \tilde{F}^{\mu\kappa} \partial_\kappa \theta + \nu^\mu,$$

Charge density

$$J_5^\mu = f^2 \partial^\mu \theta + \nu_5^\mu.$$

$\theta$  „decay constant“

Dissipative corrections  
(viscosity, resistance, etc.)

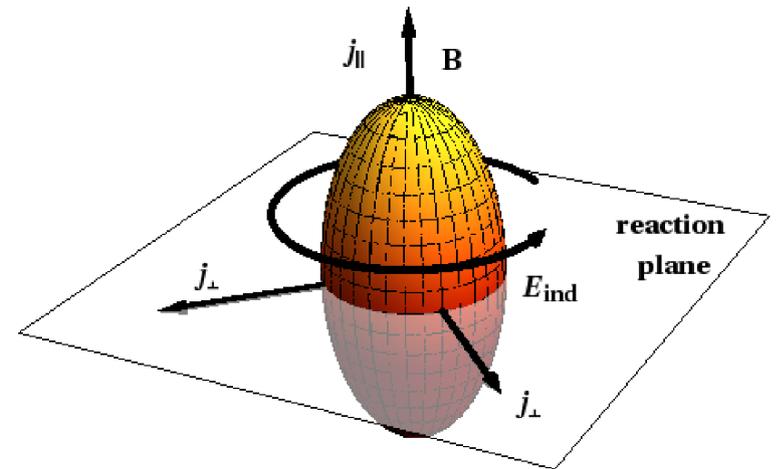
The diagram illustrates the constitutive relations for the energy-momentum tensor  $T^{\mu\nu}$ , the current  $J^\mu$ , and the axial current  $J_5^\mu$ . Arrows point from the physical quantities on the left to their corresponding terms in the equations. The energy density  $\epsilon$  and pressure  $P$  are associated with the first two terms of  $T^{\mu\nu}$ . The charge density  $\rho$  is associated with the first term of  $J^\mu$ . The decay constant  $\theta$  is associated with the first term of  $J_5^\mu$ . The terms  $\tau^{\mu\nu}$ ,  $\nu^\mu$ , and  $\nu_5^\mu$  are collectively labeled as dissipative corrections (viscosity, resistance, etc.).

**Notice the additional current**

# Phenomenology

An additional electric current induced by the  $\theta$ -field:

$$j_\lambda = -C\mu_5 B_\lambda + C\epsilon_{\lambda\alpha\kappa\beta}u^\alpha\partial^\kappa\theta E^\beta - Cu_\lambda(\partial\theta \cdot B)$$

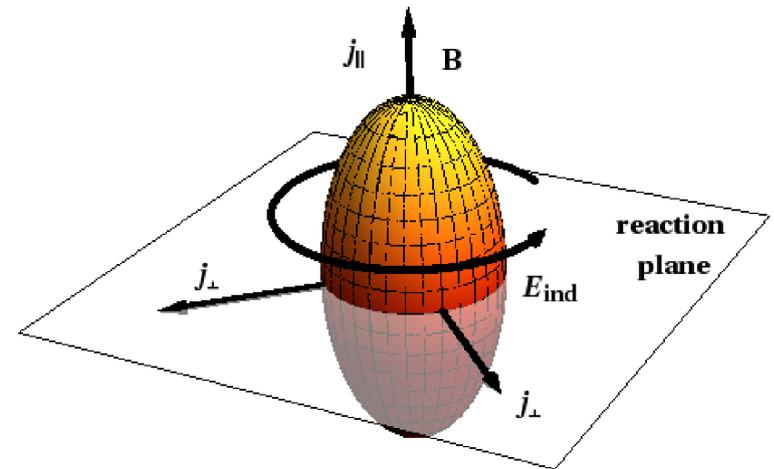


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- **Chiral Magnetic Effect** (electric current along B-field)

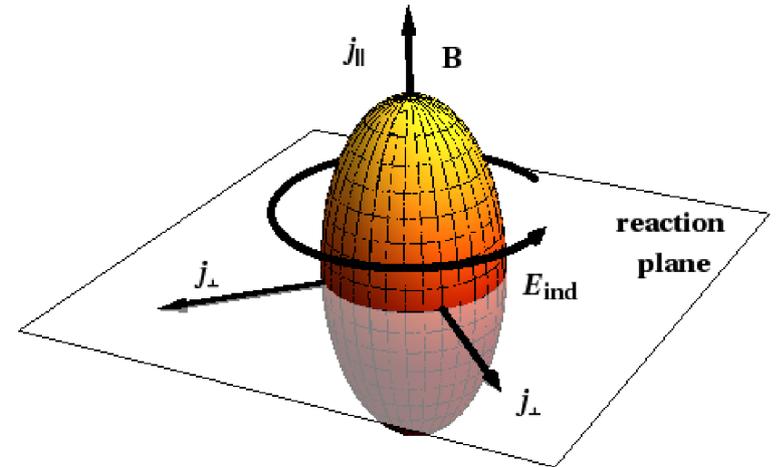


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- **Chiral Magnetic Effect** (electric current along B-field)
- **Chiral Electric Effect** (electric current transverse to E-field and to both normal and superfluid velocities)

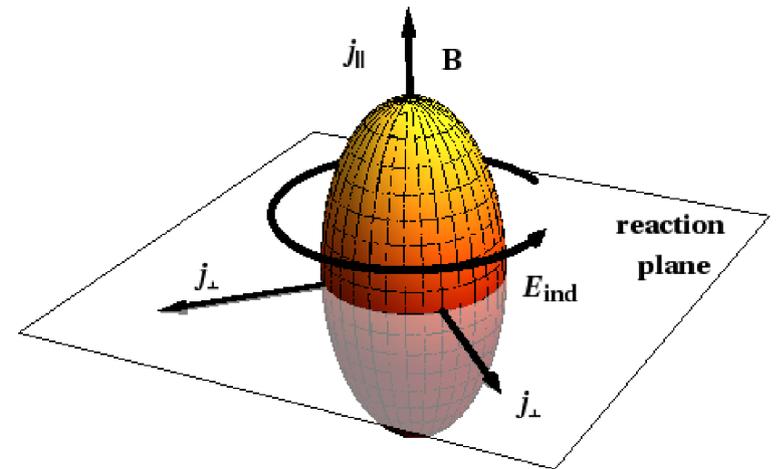


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- **Chiral Magnetic Effect** (electric current along B-field)
- **Chiral Electric Effect** (electric current transverse to E-field and to both normal and superfluid velocities)
- **Chiral Dipole Wave** (dipole moment induced by B-field)

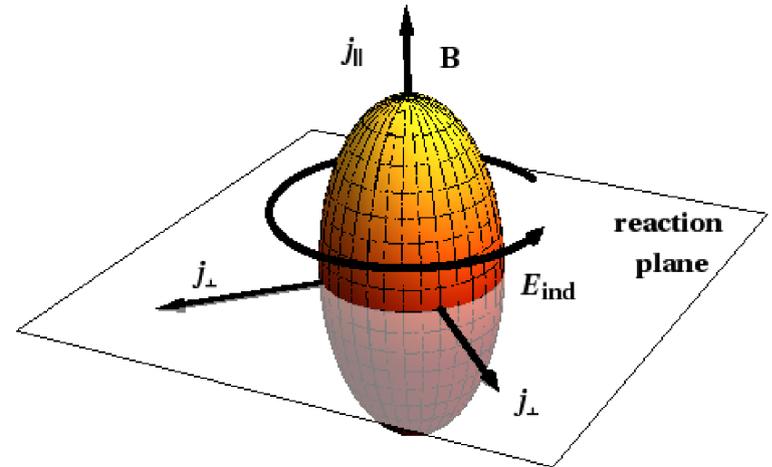


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- **Chiral Magnetic Effect** (electric current along B-field)
- **Chiral Electric Effect** (electric current transverse to E-field and to both normal and superfluid velocities)
- **Chiral Dipole Wave** (dipole moment induced by B-field)
- The field  $\theta(x)$  itself: **Chiral Magnetic Wave** (propagating imbalance between the number of left- and right-handed quarks)



# Change in entropy and higher order gradient corrections

Higher order correction obey the Landau conditions

$$u_\mu \tau^{\mu\nu} = 0, \quad u_\mu \nu^\mu = 0, \quad u_\mu \nu_5^\mu = 0$$

# Change in entropy and higher order gradient corrections

Higher order correction obey the Landau conditions

$$u_\mu \tau^{\mu\nu} = 0, \quad u_\mu \nu^\mu = 0, \quad u_\mu \nu_5^\mu = 0$$

Using both hydrodynamic equations and constitutive relations one can derive

$$\partial_\mu \left( s u^\mu - \frac{\mu}{T} \nu^\mu - \frac{\mu_5}{T} \nu_5^\mu \right) = -\frac{1}{T} (\partial_\mu u_\nu) \tau^{\mu\nu} - \nu^\mu \left( \partial_\mu \frac{\mu}{T} - \frac{1}{T} E_\mu \right) - \nu_5^\mu \partial_\mu \frac{\mu_5}{T}$$

- Entropy production is always non-negative

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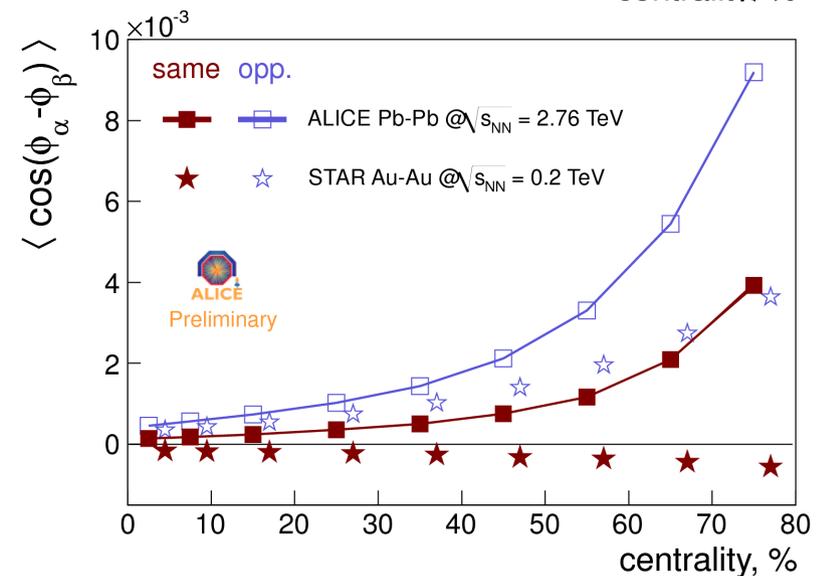
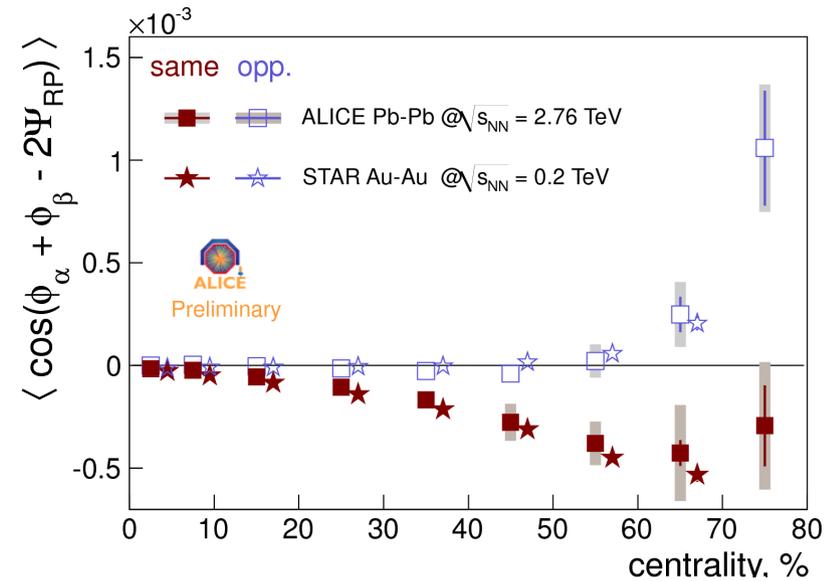
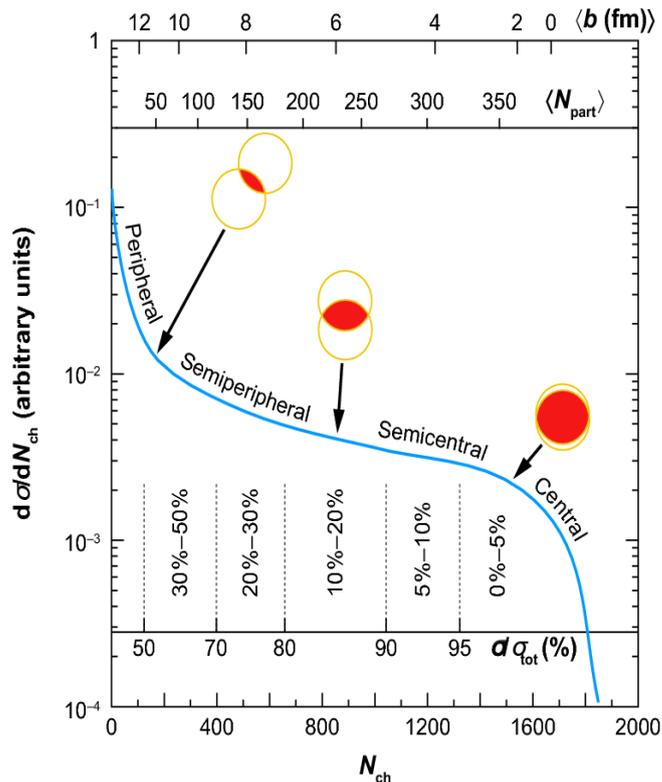
$$\partial_\mu (su^\mu - \cancel{\frac{\mu}{T}} - \cancel{\frac{\mu_5}{T}}) = -\frac{1}{T} (\cancel{\partial_\mu \nu^\mu}) \tau^{\mu\nu} - \nu^\mu (\partial_\mu \cancel{\frac{\mu}{T}} - \frac{1}{T} E_\mu) - \cancel{\nu_5^\mu \partial_\mu \frac{\mu_5}{T}}$$

- Entropy production is always non-negative
- No additional anomalous first-order corrections to the currents
- Only the “normal” component contributes to the entropy current, while the “superfluid” component has zero entropy

# Experiment

$$\gamma_{\alpha,\beta} = \langle \cos(\phi_\alpha + \phi_\beta) \rangle = \langle \cos \cos \rangle - \langle \sin \sin \rangle$$

$$\delta_{\alpha,\beta} = \langle \cos(\phi_\alpha - \phi_\beta) \rangle = \langle \cos \cos \rangle + \langle \sin \sin \rangle$$

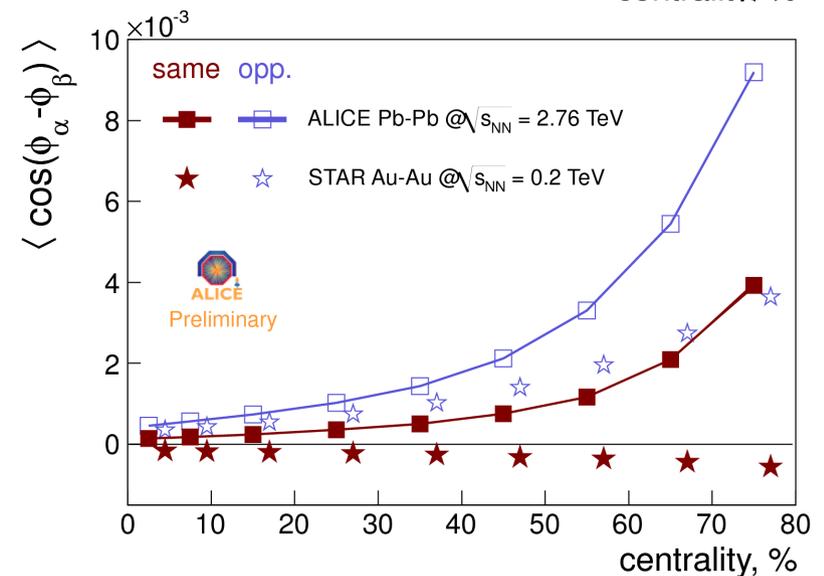
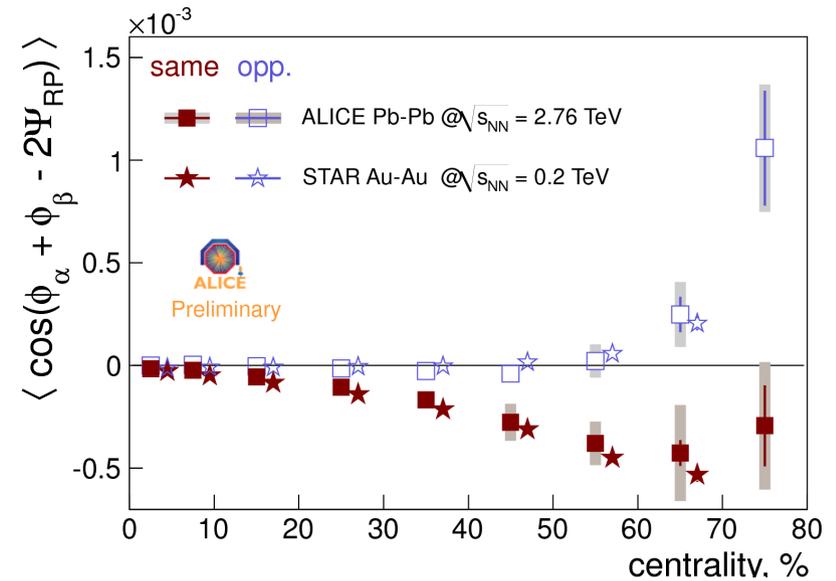


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in-plane
out-of-plane



# Experiment

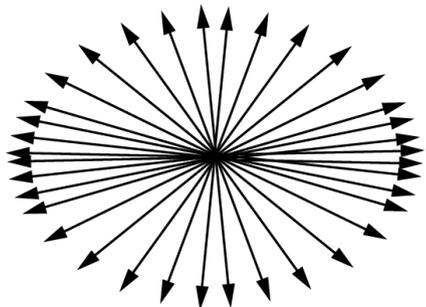
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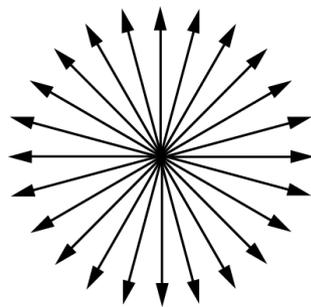
in-plane
out-of-plane

$$\gamma_{\alpha,\beta} \sim v_2 F_{\alpha,\beta} - H_{\alpha,\beta}^{\text{out}} + H_{\alpha,\beta}^{\text{in}}$$

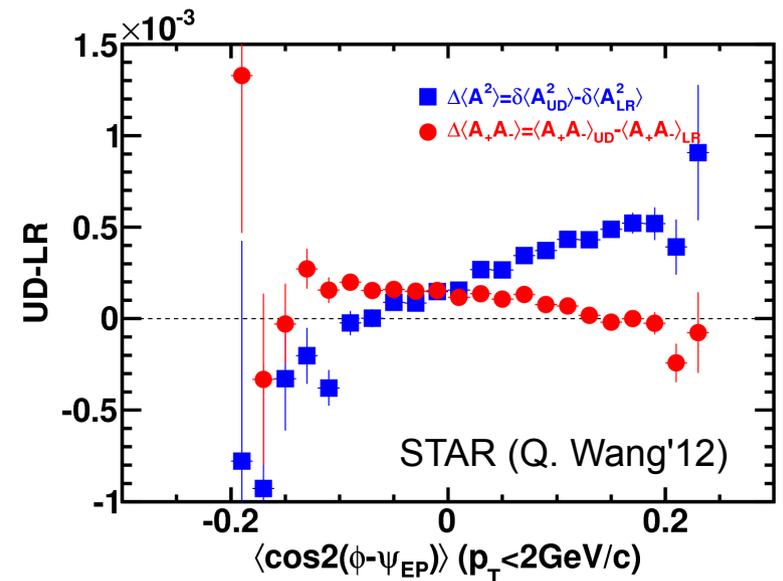
$$\delta_{\alpha,\beta} \sim F_{\alpha,\beta} + H_{\alpha,\beta}^{\text{out}} + H_{\alpha,\beta}^{\text{in}}$$



$v_2 > 0$



$v_2 = 0$



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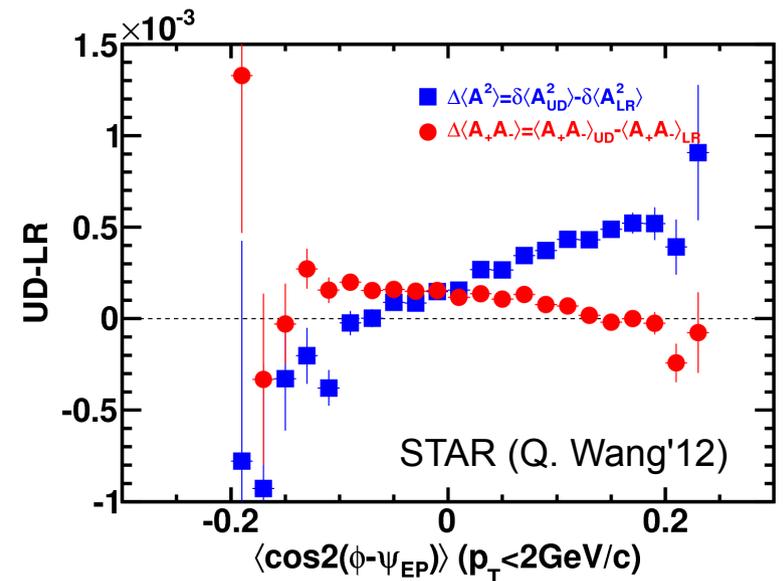
$$\delta_{\alpha,\beta} = \langle \cos(\phi_\alpha - \phi_\beta) \rangle = \langle \cos \cos \rangle + \langle \sin \sin \rangle$$

in-plane
out-of-plane

$$\gamma_{\alpha,\beta} \sim v_2 F_{\alpha,\beta} - H_{\alpha,\beta}^{\text{out}} + H_{\alpha,\beta}^{\text{in}}$$

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flow-dependent
flow-independent



# Experiment

$$\gamma_{\alpha,\beta} = \langle \cos(\phi_\alpha + \phi_\beta) \rangle = \langle \cos \cos \rangle - \langle \sin \sin \rangle$$

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in-plane
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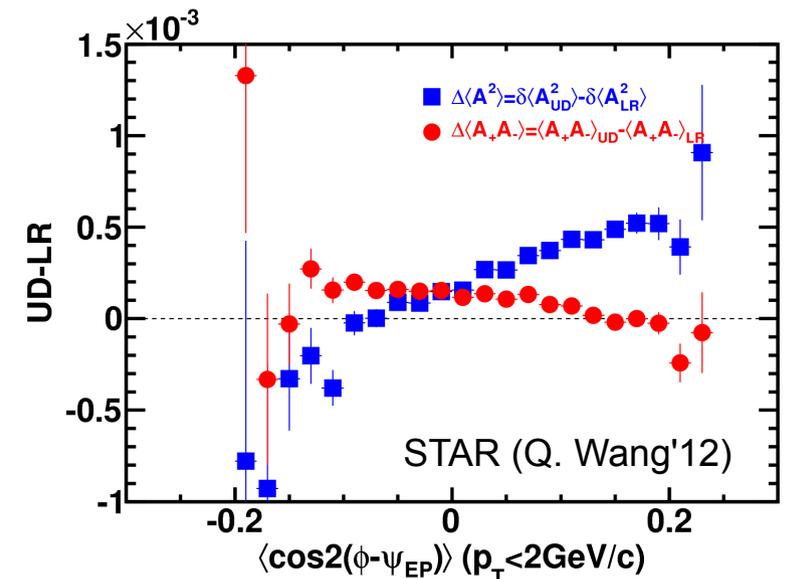
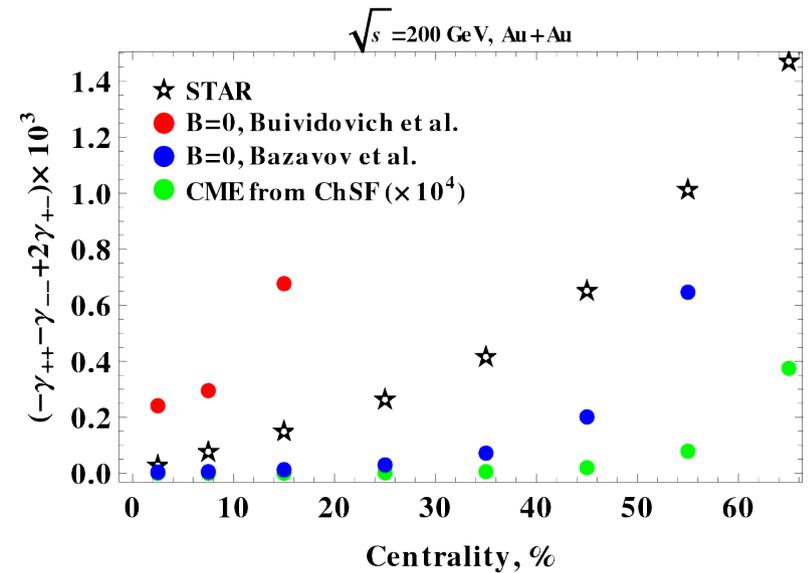
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flow-dependent
flow-independent

$$H_{++} + H_{--} - 2H_{+-} \sim \frac{4\pi\tau^2 \rho^2 \mathcal{R}^2}{3N_q^2} \left( \langle j_{\parallel}^2 \rangle + 2\langle j_{\perp}^2 \rangle \right)$$

Buividovich, Chernodub, Luschevskaya, Polikarpov' 09



# Interesting projects

- Add more flavors. The „axion-like“ field might be just a pion. At strong B it will be then a 2D pion.
- Role of the low-dimensional defects in inducing superfluidity/superconductivity. Copenhagen vacuum.
- Entropy and thermalization of various parts of fermionic spectrum. Relation to the vacuum entropy.
- Considering vortices (axionic strings). One might reproduce the chiral vortical effect.
- Holographic model of the chiral superfluidity and high-order corrections.
- Experimental searches for the mentioned effects.
- Chiral electric effect on a lattice.

**Thank you for the attention!**

**and**

**Have a good time!**

**All comments on the papers are welcome!  
Also feel free to ask questions about the experimental observables.**

# Backup slide

